LECTURE NOTES ON ENGINEERING MECHANICS

B. Tech II Semester (MR-20)

Prepared By

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MECHANICAL ENGINEERING

MALLA REDDY ENGINEERING COLLEGE (Autonomous)

ENGINEERING MECHANICS

| | Category | Hours / Week | | | Credits | Maximum Marks | | | |
|--|--|--|---|--|---|---|---|---|--|
| A0303 | Coro | L | Т | Р | С | CIA | SEE | Total | |
| | Core | 3 | 0 | 0 | 3 | 30 | 70 | 100 | |
| Contact Classes: 48 | Tutorial Classes: Nil | Practical Classes: Nil | | | | | Total Classes: 48 | | |
| COURSE OBJECTIV The course should enab I. Students should required for analy II. Identify an appro- environment, mo III. Understand the r methods and met IV. To solve the proble and vibrations fo Mechanics of Flu | ES: le the students to: develop the ability to work yzing static structures. opriate structural system to del the problem using good neaning of centre of gravity hod of moments m of equilibrium by using th r preparing the students for l uids, Mechanical Design and | t comf study free-b (mas (mas e prin nigher Struc | fortably ing a g ody dia ss)/cen iciple o level o tural A | y with given pr agrams troid an f work courses nalysis | basic engin roblem and and accura id moment and energy such as Me etc | eering n l isolate te equilib of Inerti , impulse echanics | nechanic it from i prium equal a using i e momen of Solids | s concept ts uations. integratio tum | |
| MODULE-I Introduction to Mechanics & System of Forces | | | | | | | Class | Classes: 10 | |
| Cypes of Forces - Con Polygon and Parallelog Couples - Free Body D Equilibrium, Equations | accurrent and non-concurren gram Law of Forces - Mor iagrams, Types of Supports of Equilibrium, Conditions | t Fore nent of and th of Eq | ces - (of Ford leir rea uilibriu | Compositions Compo | ition of fo its Applica Internal and mi's Theore | rce – Re tion - V d Externa em. | esultant arignon' Il Forces | - Triangl s theorem - Types o | |
| MODULE-II Fri | on, Centroid and Center of Gravity | | | | | | Class | ses: 10 | |
| Friction: Types of fric aws of friction. Motion Centroids of Lines and | tion, Limiting friction, Law n of bodies - wedge, screw, a Areas - simple figures - Ce c, composite solids - Centroid | rs of t screw entroid ds of y | friction jack. (d of co volume | , static Centroio mposito s. | and dynan 1 and Cente e figures. P | nic fricti er of Gra appus th | on, appl vity: Intr eorem - | ication of oduction Centre of | |
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| gravity of simple solids | oment of Inertia | | | | | | Class | ses: 10 | |
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MODULE-V Work, Power, Energy & Mechanical Vibrations Classes: 09

Work, Power and Energy: Introduction, work-energy equation - motion of connected bodies - work done by a spring - general plane motion. Mechanical Vibrations: Definitions, concepts - simple harmonic motion - free vibrations - Simple and compound pendulums.

Text Books:

- 1.S. Timoshenko, D.H. Young, J.V. Rao and Sukumar Pati, "Engineering Mechanics", Tata McGraw-Hill Education, 5th Edition, 2013.
- 2.K.Vijaya Kumar Reddy, J. Suresh Kumar, "Engineering Mechanics", B S Publications, 3 rd Edition, 2013.

Reference Books:

- 1. Beer, F.P and Johnston Jr. E.R. "Vector Mechanics for Engineers", Tata McGraw-Hill Education 10th Edition (India) Pvt Ltd.. 2013.
- 2. Fedinand. L. Singer, "Engineering Mechanics", Harper & Row Publishers, 3rd Edition, 1975.
- 3. S.Bhavikatti, "ATextBookofEngineeringMechanics", NewAgeInternational, 1st Edition, 2012
- 4. R.S. Khurmi, "A Text Book of Engineering Mechanics", S.Chand Publications, 21st Edition, 2007.
- 5. K L Kumar, "Engineering Mechanics", Tata McGraw Hill Education, 4th Edition, 2011

Web References:

- 1. http://nptel.ac.in/courses/112103109/
- 2. http://nptel.ac.in/courses/112106180/
- 3. http://nptel.ac.in/courses/115104094/

E-Text Books:

- 1. http://www.mathalino.com/reviewer/engineering-mechanics/equilibrium-force-system
- 2. http://ascelibrary.org/journal/jenmdt
- 3. <u>https://tll.mit.edu/sites/default/files/SUTDVideoThumb/freebodydiagrams.pdf</u>

MODULE I

Introduction to Mechanics & System of Forces

Mechanics

It is defined as that branch of science, which describes and predicts the conditions of rest or motion of bodies under the action of forces. Engineering mechanics applies the principle of mechanics to design, taking into account the effects of forces.

Statics

Statics deal with the condition of equilibrium of bodies acted upon by forces.

<u>Rigid body</u>

A rigid body is defined as a definite quantity of matter, the parts of which are fixed in position relative to each other. Physical bodies are never absolutely but deform slightly under the action of loads. If the deformation is negligible as compared to its size, the body is termed as rigid.



<u>Force</u>

Force may be defined as any action that tends to change the state of rest or motion of a body to which it is applied.

The three quantities required to completely define force are called its specification or characteristics. So the characteristics of a force are:

- 1. Magnitude
- 2. Point of application
- 3. Direction of application



Concentrated force/point load



Distributed force



Line of action of force

The direction of a force is the direction, along a straight line through its point of application in which the force tends to move a body when it is applied. This line is called line of action of force.

Representation of force

Graphically a force may be represented by the segment of a straight line.



Composition of two forces

The reduction of a given system of forces to the simplest system that will be its equivalent is called the problem of composition of forces.

Parallelogram law

If two forces represented by vectors AB and AC acting under an angle α are applied to a body at point A. Their action is equivalent to the action of one force, represented by vector AD, obtained as the diagonal of the parallelogram constructed on the vectors AB and AC directed as shown in the figure.



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1

Force AD is called the resultant of AB and AC and the forces are called its components.



$$R = \sqrt{\left(P^2 + Q^2 + 2PQ \times \cos\alpha\right)}$$

Now applying triangle law

$$\frac{P}{Sin\gamma} = \frac{Q}{Sin\beta} = \frac{R}{Sin(\pi - \alpha)}$$

Special cases

Case-I: If
$$\alpha = 0^{\circ}$$

 $R = \sqrt{\left(P^2 + Q^2 + 2PQ \times Cos0^2\right)} = \sqrt{\left(P + Q\right)^2} = (P + Q)$
 $\stackrel{P}{\longrightarrow} Q \qquad R$
 $R = P + Q$

Case- II: If $\alpha = 180^{\circ}$

$$R = \sqrt{(P^2 + Q^2 + 2PQ \times Cos180^2)} = \sqrt{(P^2 + Q^2 - 2PQ)} = \sqrt{(P - Q)^2} = (P - Q)$$
Q



Resolution of a force

The replacement of a single force by a several components which will be equivalent in action to the given force is called resolution of a force.



Action and reaction

Often bodies in equilibrium are constrained to investigate the conditions.



Free body diagram

Free body diagram is necessary to investigate the condition of equilibrium of a body or system. While drawing the free body diagram all the supports of the body are removed and replaced with the reaction forces acting on it.

- 1. Draw the free body diagrams of the following figures.

2. Draw the free body diagram of the body, the string CD and thering.





3. Draw the free body diagram of the following figures.



Equilibrium of colinear forces:

Equilibrium law: Two forces can be in equilibrium only if they are equal in magnitude, opposite in direction and collinear in action.



Superposition and transmissibility

Problem 1: A man of weight W = 712 N holds one end of a rope that passes over a pulley vertically above his head and to the other end of which is attached a weight Q = 534 N. Find the force with which the man's feet press against the floor.

Tension in the string S is equal
to the load attached tait
$$R = 534N$$
,
So $S = 5324N$.
Now applying
porollalogram law
resultant force
 $R = \sqrt{W^2 + 4^2 + 2WSCOS/S0^{1}}$
 $= \sqrt{W^2 + 5^2 - 2WS}$
 $= \sqrt{W^2 + 5^2 - 2WS}$
 $= \sqrt{W^2 + 5^2 - 2WS}$
 $= \sqrt{W^2 - 534 \subseteq 178N(4)}$
Readtion on the man's feet $= 178N(4)$

Problem 2: A boat is moved uniformly along a canal by two horses pulling with forces P = 890 N and Q = 1068 N acting under an angle $\alpha = 60^{\circ}$. Determine the magnitude of the resultant pull on the boat and the angles β and ν .





Resolution of a force

Replacement of a single force by several components which will be equivalent in action to the given force is called the problem of resolution of aforce.

By using parallelogram law, a single force R can be resolved into two components P and Q intersecting at a point on its line of action.



Equilibrium of collinear forces:

Equilibrium law: Two forces can be in equilibrium only if they are equal in magnitude, opposite in direction and collinear in action.



Law of superposition

The action of a given system of forces on a rigid body will no way be changed if we add to or subtract from them another system of forces in equilibrium.

Problem 3: Two spheres of weight P and Q rest inside a hollow cylinder which is resting on a horizontal force. Draw the free body diagram of both the spheres, together and separately.



Problem 4: Draw the free body diagram of the figure shown below.











Problem 6: Find the reactions R_1 and R_2 .



Problem 7: Two rollers of weight P and Q are supported by an inclined plane and vertical walls as shown in the figure. Draw the free body diagram of both the rollers separately.



Problem 8: Find θ_n and θ_t in the following figure.



Problem 9: For the particular position shown in the figure, the connecting rod BA of an engine exert a force of P = 2225 N on the crank pin at A. Resolve this force into two rectangularcomponents P_h and P_v horizontally and vertically respectively at A.



 $P_h = 2081.4 \text{ N}$ $P_v = 786.5 \text{ N}$

Equilibrium of concurrent forces in a plane

- •If a body known to be in equilibrium is acted upon by several concurrent, coplanar forces, then these forces or rather their free vectors, when geometrically added must form a closedpolygon.
- •This system represents the condition of equilibrium for any system of concurrent forces in aplane.







Lami's theorem

If three concurrent forces are acting on a body kept in an equilibrium, then each force is proportional to the sine of angle between the other two forces and the constant of proportionality issame.



$$\frac{P}{\sin\alpha} = \frac{Q}{\sin\beta} = \frac{R}{\sin U}$$



Problem: A ball of weight Q = 53.4N rest in a right angled trough as shown in figure. Determine the forces exerted on the sides of the trough at D and E if all the surfaces are perfectlysmooth.



Problem: An electric light fixture of weight Q = 178 N is supported as shown in figure. Determine the tensile forces S_1 and S_2 in the wires BA and BC, if their angles of inclination are given.









 $S_1 \cos \alpha = P$

 $S = Psec\alpha$

$$R_b = W + S \sin \alpha$$
$$= W + \frac{P}{\cos \alpha} \times \sin \alpha$$
$$= W + P \tan \alpha$$

Problem: A right circular roller of weight W rests on a smooth horizontal plane and is held in position by an inclined bar AC. Find the tensions in the bar AC and vertical reaction R_b if there is also a horizontal force P is active.



Theory of transmissiibility of a force:

The point of application of a force may be transmitted along its line of action without changing the effect of force on any rigid body to which it may be applied.

Problem:





$$\sum_{S_1 \text{ cos } 30 + 20 \text{ sin } 60 = S_2 \text{ sin } 30}$$
$$\frac{\sqrt{3}}{2} \sum_{1}^{S_1 + 20} \frac{\sqrt{3}}{2} = \frac{S_2}{2}$$
$$\frac{S_2}{2} = \frac{\sqrt{3}}{2} \sum_{1}^{S_1 + 10} \sqrt[3]{-1}$$
$$S_2 = \sqrt{3}S_1 + 20\sqrt{3}$$

$$\sum Y = 0$$

 $S_1 \sin 30 + S_2 \cos 30 = S_a \cos 60 + 20$
 $\frac{S_1}{2} + S_2 \frac{\sqrt{3}}{2} = \frac{20}{2} + 20$
 $\frac{S_1}{2} + \frac{\sqrt{3}}{2} S = 30$
 $S_1 + \sqrt{3}S_2 = 60$

Substituting the value of S_2 in Eq.2, we get

$$S_{1} + \sqrt{3} (\sqrt[3]{S_{1}} + 20 \sqrt[3]{}) = 60$$

$$S_{1} + 3S_{1} + 60 = 60$$

$$4S_{1} = 0$$

$$S_{1} = 0KN$$

$$S_{2} = 20 \sqrt{3} = 34.64KN$$

(1)

(2)

Problem: A ball of weight W is suspended from a string of length 1 and is pulled by a horizontal force Q. The weight is displaced by a distance d from the vertical position as shown in Figure. Determine the angle α , forces Q and tension in the string S in the displaced position.









$$= \sqrt{1 - \frac{d^2}{l^2}}$$
$$= \frac{1}{l} l \sqrt{1 - d^2}$$

Applying Lami's theorem,

$$\frac{S}{\sin 90} = \frac{Q}{\sin(90+\alpha)} = \frac{W}{\sin(180-\alpha)}$$

$$\frac{Q}{\sin(90+\alpha)} = \frac{W}{\sin(180-\alpha)}$$
$$\Rightarrow Q = \frac{W\cos\alpha}{\sin\alpha} = \frac{W \int_{l^+}^{l^+} \frac{1}{l} \sqrt{l^2-d^2}}$$
$$\Rightarrow Q = \frac{Wd}{\sqrt{l^2-d^2}}$$
$$S = \frac{W}{\sin\alpha} = \frac{W}{\frac{1}{l} \sqrt{l^2-d^2}}$$
$$= \frac{Wl}{\sqrt{l^2-d^2}}$$

Problem: Two smooth circular cylinders each of weight W = 445 N and radius r = 152 mm are connected at their centres by a string AB of length l = 406 mm and rest upon a horizontal plane, supporting above them a third cylinder of weight Q = 890 N and radius r = 152 mm. Find the forces in the string and the pressures produced on the floor at the point of contact.



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Problem: Two identical rollers each of weight Q = 445 N are supported by an inclined plane and a vertical wall as shown in the figure. Assuming smooth surfaces, find the reactions induced at the points of support A, B and C.



 $\frac{R_a}{\sin 120} \quad \frac{S}{\sin 150} = \frac{445}{\sin 90}$

 $\Rightarrow R_a = 385.38N$ $\Rightarrow S = 222.5N$

S = 399.27*N*



Problem:

A weight Q is suspended from a small ring C supported by two cords AC and BC. The cord AC is fastened at A while cord BC passes over a frictionless pulley at B and carries a weight P. If P = Q and $\alpha = 50^{\circ}$, find the value of β .



Resolving horizontally

 $\sum_{S \le 0} X = 0$ S \sin 50 = Q \sin \beta Resolving vertically

 $\sum_{S \to S} Y = 0$ $\Rightarrow S \cos 50 + Q \sin \beta = Q$ $\Rightarrow S \cos 50 = Q(1 - \cos \beta)$ Putting the value of S from Eq. 1, we get (1)

R_c

 $S \cos 50 + Q \sin \beta = Q$ $\Rightarrow S \cos 50 = Q(1 - \cos\beta)$ $\Rightarrow Q \frac{\sin\beta}{\sin 50} \cos 50 = Q(1 - \cos\beta)$ $\Rightarrow \cot 50 = \frac{1 - \cos\beta}{\sin\beta}$ $\Rightarrow 0.839 \sin\beta = 1 - \cos\beta$

Squaring both sides, $0.703\sin^2\beta = 1+\cos^2\beta - 2\cos\beta$ $0.703(1-\cos^2\beta) = 1+\cos^2\beta - 2\cos\beta$ $0.703 - 0.703\cos^2\beta = 1+\cos^2\beta - 2\cos\beta$ $\Rightarrow 1.703\cos^2\beta - 2\cos\beta + 0.297 = 0$ $\Rightarrow\cos^2\beta - 1.174\cos\beta + 0.297 = 0$ $\Rightarrow\beta = 63.13^{\Box}$

Method of moments

Moment of a force with respect to a point:



- •Considering wrench subjected to two forces P and Q of equal magnitude. It is evident that force P will be more effective compared to Q, though they are of equalmagnitude.
- •The effectiveness of the force as regards it is the tendency to produce rotation of a body about a fixed point is called the moment of the force with respect to that point.
- •Moment = Magnitude of the force × Perpendicular distance of the line of action offorce.
- •Point O is called moment centre and the perpendicular distance (i.e. OD) is called momentarm.
- •Unit isN.m

Theorem of Varignon:

The moment of the resultant of two concurrent forces with respect to a centre in their plane is equal to the alzebric sum of the moments of the components with respect to some centre.

Problem 1:

A prismatic clear of AB of length l is hinged at A and supported at B. Neglecting friction, determine the reaction R_b produced at B owing to the weight Q of the bar.

Taking moment about point A,

$$R_{b} \times l = Q \cos \alpha . \frac{l}{2}$$

 $\Rightarrow R_{b} = \frac{Q}{\cos \alpha 2}$



Problem 2:

A bar AB of weight Q and length 2l rests on a very small friction less roller at D and against a smooth vertical wall at A. Find the angle α that the bar must make with the horizontal in equilibrium.



Resolving vertically, $R_d \cos \alpha = Q$

Now taking moment about A,

$$\frac{R_{d.a}}{\cos \alpha} - Q.l\cos \alpha = 0$$

$$\Rightarrow \frac{Q.a}{\cos^2 \alpha} - Q.l\cos \alpha = 0$$

$$\Rightarrow Q.a - Q.l\cos^3 \alpha = 0$$

$$\Rightarrow \cos^3 \alpha = \frac{Q.a}{Q.l}$$

$$\Rightarrow \alpha = \cos^{-1} \sqrt[3]{\frac{a}{l}}$$

Problem 3:

If the piston of the engine has a diameter of 101.6 mm and the gas pressure in the cylinder is 0.69 MPa. Calculate the turning moment M exerted on the crankshaft for the particular configuration.



Area of cylinder $A = \frac{\pi}{4} (0.1016)^2 = 8.107 \times 10^{-3} m^2$

Force exerted on connectingrod,

$$F = Pressure \times Area = 0.69 \times 10^{6} \times 8.107 \times 10^{-3} = 5593.83 N$$

Now
$$\alpha = \sin^{-1} \left(\frac{178}{380} \right)^{-1} = 27.93^{-1}$$

 $S \cos \alpha = F$ $\Rightarrow S = \frac{F}{\cos \alpha} = 6331.29N$

Now moment entered on crankshaft,

 $S\cos\alpha \times 0.178 = 995.7N = 1KN$



Problem 4:

A rigid bar AB is supported in a vertical plane and carrying a load Qat its free end. Neglecting the weight of bar, find the magnitude of tensile force S in the horizontal string CD.



Taking moment about A,

$$\sum M_{A} = 0$$

S. $\frac{l}{\cos \alpha} = Q.l\sin \alpha 2$
 $\Rightarrow S = \frac{Q.l\sin \alpha}{l}$
 $\Rightarrow S = 2Q.\tan \alpha$

MODULE II

Friction, Centroid and Center of Gravity

Friction

- •The force which opposes the movement or the tendency of movement is called **Frictional force or simply friction**. It is due to the resistance to motion offered by minutely projecting particles at the contact surfaces. However, there is a limit beyond which the magnitude of this force cannotincrease.
- •If the applied force is more than this limit, there will be movement of one body over the other. This limiting value of frictional force when the motion is impending, it is known as **LimitingFriction**.
- •When the applied force is less than the limiting friction, the body remains at rest and such frictional force is called **Static Friction**, which will be having any value between zero and the limitingfriction.
- •If the value of applied force exceeds the limiting friction, the body starts moving over the other body and the frictional resistance experienced by the body while moving is known as **Dynamic Friction**. Dynamic friction is less than limitingfriction.
- •Dynamic friction is classified into following twotypes:
- a) Slidingfriction
- b) Rolling friction
- •Sliding friction is the friction experienced by a body when it slides over the other body.
- •Rolling friction is the friction experienced by a body when it rolls over a surface.
- •It is experimentally found that the magnitude of limiting friction bears a constant ratio to the normal reaction between two surfaces and this ratio is called **Coefficientof Friction**.



$$\overline{N}$$

where F is limiting friction and N is normal reaction between the contact surfaces.

Coefficient of friction is denoted by μ .

Thus,
$$\mu = \frac{F}{N}$$

Laws of friction

- 1. The force of friction always acts in a direction opposite to that in which body tends tomove.
- 2. Till the limiting value is reached, the magnitude of friction is exactly equal to the force which tends to move thebody.
- 3. The magnitude of the limiting friction bears a constant ratio to the normal reaction between the two surfaces of contact and this ratio is called coefficient offriction.
- 4. The force of friction depends upon the roughness/smoothness of thesurfaces.
- 5. The force of friction is independent of the area of contact between the two surfaces.
- 6. After the body starts moving, the dynamic friction comes into play, the magnitude of which is less than that of limiting friction and it bears a constant ratio with normal force. This ratio is called **coefficient of dynamicfriction**.

Angle of friction

Consider the block shown in figure resting on a horizontal surface and subjected to horizontal pull P. Let F be the frictional force developed and N the normal reaction. Thus, at contact surface the reactions are F and N. They can be graphically combined to get the reaction R which acts at angle θ to normal reaction. This angle θ called the angle of friction is given by

$$\tan \theta = \frac{F_{---}}{N}$$

As P increases, F increases and hence θ also increases. θ can reach the maximum value α when F reaches limiting value. At this stage,

This value of α is called Angle of Limiting Friction. Hence, the angle of limiting friction may be defined as the angle between the resultant reaction and the normal to the plane on which the motion of the body is impending.

Angle of repose



Consider the block of weight W resting on an inclined plane which makes an angle θ with the horizontal. When θ is small, the block will rest on the plane. If θ is gradually increased, a stage is reached at which the block start sliding down the plane. The angle θ for which the motion is impending, is called the angle of repose. Thus, the maximum inclination of the plane on which a body, free from external forces, can repose is called **Angle of Repose**.

Resolving vertically, N = W. $\cos \theta$

Resolving horizontally, $F = W. \sin \theta$

Thus, $\tan \theta = F$

If ϕ is the value of θ when the motion is impending, the frictional force will be limiting friction and hence,

$$\tan \varphi = \frac{F}{N} = \frac{1}{N}$$
$$= \mu = \tan \alpha$$
$$\Rightarrow \varphi = \alpha$$
Thus, the val

Thus, the value of angle of repose is same as the value of limiting angle of repose.

Cone of friction



- •When a body is having impending motion in the direction of force P, the frictional force will be limiting friction and the resultant reaction R will make limiting angle α with thenormal.
- •If the body is having impending motion in some other direction, the resultant reaction makes limiting frictional angle α with the normal to that direction. Thus, when the direction of force P is gradually changed through 360°, the resultant R generates a right circular cone with semi-central angle equal to α .

Problem 1: Block A weighing 1000N rests over block B which weighs 2000N as shown in figure. Block A is tied to wall with a horizontal string. If the coefficient of friction between blocks A and B is 0.25 and between B and floor is 1/3, what should be the value of P to move the block (B), if

- (a) P ishorizontal.
- (b) P acts at 30° upwards tohorizontal.

Solution: (a)



Considering block A,

$$\sum V = 0$$
$$N_1 = 1000N$$

Since F₁ is limiting friction,

$$\frac{F_1}{N_1} = \mu = 0.25$$

$$F_1 = 0.25N_1 = 0.25 \times 1000 = 250N$$

 $\sum H = 0$ $F_1 - T = 0$ $T = F_1 = 250N$

Considering equilibrium of block B,

$$\sum_{N_2 = 0}^{N_2 = 0} V = 0$$

N₂ = 2000 -N₁ = 0
N₂ = 2000 +N₁ = 2000 +1000 = 3000N

$$\frac{F_2}{N_2} = \mu = \frac{1}{3}$$

 $F_2 = 0.3N_2 = 0.3 \times 1000 = 1000N$

$$\sum H = 0$$

P = F₁ + F₂ = 250 + 1000 = 1250N

(b) When P is inclined:

$$\sum V = 0$$

 $N_2 - 2000 - N_1 + P.\sin 30 = 0$
 $\Rightarrow N_2 + 0.5P = 2000 + 1000$
 $\Rightarrow N_2 = 3000 - 0.5P$

From law of friction,

$$F = \frac{1}{3} \frac{N}{2} = \frac{1}{3} (3000 - 0.5P) = 1000 - \frac{0.5P}{3}$$

$$\sum H = 0$$

$$P \cos 30 = F_1 + F_2 = 0.5$$

$$\Rightarrow P \cos 30 = 250 + 1000 - P$$

$$= 1250$$

$$\Rightarrow P \left(\cos 30 + \frac{0.5P}{3} \right) = 1250$$

$$\Rightarrow P = 1210.43N$$

Problem 2: A block weighing 500N just starts moving down a rough inclined plane when supported by a force of 200N acting parallel to the plane in upward direction. The same block is on the verge of moving up the plane when pulled by a force of 300N acting parallel to the plane. Find the inclination of the plane and coefficient of friction between the inclined plane and theblock.



$$\sum_{N=500.\cos\theta} V = 0$$

F₁= $\mu N = \mu.500 \cos\theta$

 $\sum H = 0$ 200 +F₁ = 500.sin θ \Rightarrow 200 + μ .500 cos θ = 500.sin θ

 $\sum_{N=500.\cos\theta} V = 0$ $F_2 = \mu N = \mu.500.\cos\theta$

 $\sum H = 0$ 500 sin θ +F₂ = 300 \Rightarrow 500 sin θ + μ .500 cos θ = 300 Adding Eqs. (1) and (2), we get

$$500 = 1000. \sin\theta$$

 $\sin \theta = 0.5$
 $\theta = 30^{\circ}$

Substituting the value of θ in Eq. 2, 500 sin 30 + μ .500 cos 30 = 300

$$\mu = \frac{50}{500 \cos 30} = 0.11547$$



(1)

Parallel forces on a plane

Like parallel forces: Coplanar parallel forces when act in the same direction. Unlike parallel forces: Coplanar parallel forces when act in different direction. Resultant of

like parallel forces:

Let P and Q are two like parallel forces act at points A and B. R = P + Q

Resultant of unlikeparallelforces: R = P-Q

R is in the direction of the force havinggreatermagnitude.

Couple:

Two unlike equal parallel forces form a couple.



The rotational effect of a couple is measured by its moment.

Moment = $P \times 1$

Sign convention: Anticlockwise couple (Positive) Clockwise couple (Negative)



Problem 1 :A rigid bar CABD supported as shown in figure is acted upon by two equal horizontal forces P applied at C and D. Calculate the reactions that will be induced at the points of support. Assume l = 1.2 m, a = 0.9 m, b = 0.6 m.



Taking moment about A, $R_a = R_b$ $R_b \times l + P \times b = P \times a$ $\Rightarrow R_b = \frac{P(0.9 - 0.6)}{1.2}$ $\Rightarrow R_b = 0.25P(\uparrow)$ $\Rightarrow R_a = 0.25P(\downarrow)$

Problem 2: Owing to weight W of the locomotive shown in figure, the reactions at the two points of support A and B will each be equal to W/2. When the locomotive is pulling the train and the drawbar pull P is just equal to the total friction at the points of contact A and B, determine the magnitudes of the vertical reactions R_a and R_b .



Taking moment about B,

$$\sum M_{B} = 0$$

$$R_{a} \times 2a + P \times b = W \times a$$

$$\Rightarrow R_{a} = \frac{W \cdot a - P \cdot b}{2a}$$

$$\therefore R_{b} = W - R_{a}$$

$$\Rightarrow R = W - \frac{W \cdot a - P \cdot b}{2a}$$

$$\Rightarrow R_{b} = \frac{W \cdot a + P \cdot b}{2a}$$

Problem 3: The four wheels of a locomotive produce vertical forces on the horizontal girder AB. Determine the reactions R_a and R_b at the supports if the loads P = 90 KN each and Q = 72 KN (All dimensions are in m).



Problem 4: The beam AB in figure is hinged at A and supported at B by a vertical cord which passes over a frictionless pulley at C and carries at its end a load P. Determine the distance x from A at which a load Q must be placed on the beam if it is to remain in equilibrium in a horizontal position. Neglect the weight of thebeam.




Problem 5: A prismatic bar AB of weight Q = 44.5 N is supported by two vertical wires at its ends and carries at D a load P = 89 N as shown in figure. Determine the forces S_a and S_b in the two wires.



Resolving vertically,

$$\sum_{a} V = 0$$

$$\Rightarrow S_a + S_b = P + Q$$

$$\Rightarrow S_a + S_b = 89 + 44.5$$

$$\Rightarrow S_a + S_b = 133.5N$$



$$\sum_{b} M_{A} = 0$$

$$\sum_{b} N_{A} = P \times \frac{l}{4} + Q \times \frac{l}{2}$$

$$\Rightarrow S_{b} = \frac{P}{4} + \frac{Q}{2}$$

$$\Rightarrow S_{b} = \frac{89}{4} + \frac{44.5}{2}$$

$$\Rightarrow S_{b} = 44.5$$

$$\therefore S_{a} = 133.5 - 44.5$$

$$\Rightarrow S_{a} = 89N$$

Centre of gravity

Centre of gravity: It is that point through which the resultant of the distributed gravity force passes regardless of the orientation of the body in space.

As the point through which resultant of force of gravity (weight) of the bodyacts.

Centroid: Centroid of an area lies on the axis of symmetry if it exits.

Centre of gravity is applied to bodies with mass and weight and centroid is applied to plane areas.



$$x=y_{c} = \frac{\text{Moment of area}}{\text{Totalarea}}$$
$$x=\int \frac{x.dA}{A}$$
$$y=\int \frac{y.dA}{A}$$

Problem 1: Consider the triangle ABC of base 'b' and height 'h'. Determine the distance of centroid from the base.



Let us consider an elemental strip of width ' b_1 ' and thickness 'dy'.

$$\Delta AEF \Box \Delta ABC$$

$$\therefore \frac{b_1 - h - y b}{h}$$

$$\Rightarrow b = b \frac{(h - y)}{h}$$

$$\Rightarrow b = b \frac{(1 - y)}{h}$$

$$\Rightarrow b = b \frac{(1 - y)}{h}$$

Area of element EF (dA) = b_{11}

$$= b$$

Area of element EF (dA) =
$$b_1 \not\land dy_y$$

= $b 1 - dy$
 $\begin{vmatrix} -h \end{vmatrix}$

$$y = \frac{y \cdot dA}{h \cdot A} (y)$$

$$= \frac{\int y \cdot b \left(1 - \frac{y}{h}\right)}{\frac{1}{2} \cdot b \cdot h}$$

$$= \frac{\int y^2}{\frac{y^2}{2} - \frac{y^3}{3h}} = \frac{1}{\frac{b \cdot h}{2}}$$

$$= \frac{2 \left[h^2 - h^3\right]}{h \left[2 - \frac{y^3}{3}\right]}$$

$$= \frac{2 \left[h^2 - h^3\right]}{h \left[2 - \frac{y^3}{3}\right]}$$

$$= \frac{2 \left[h^2 - h^3\right]}{h \left[2 - \frac{y^3}{3}\right]}$$

3Therefore, y_c is at a distance of h/3 from base.

Problem 2: Consider a semi-circle of radius R. Determine its distance from diametral axis.



Due to symmetry, centroid ' y_c ' must lie on Y-axis.

Consider an element at a distance 'r' from centre 'o' of the semicircle with radial width dr.

Area of element = $(r.d\theta) \times dr$

Moment of area about
$$x = \int y.dA$$

$$= \int_{0}^{\pi R} (r.d\theta).dr \times (r.\sin\theta)$$

$$= \int_{0}^{\pi R} \int_{0}^{r^{2}} \sin\theta.dr.d\theta$$

$$= \int_{0}^{\pi R} (r^{2}.dr).\sin\theta.d\theta$$

$$= \int_{0}^{\pi R^{3}} \int_{0}^{R} \sin\theta.d\theta$$

$$= \int_{0}^{R} \int_{0}^{r} -\cos\theta \int_{0}^{\pi} \int_{0}^{r} \frac{R^{3}}{3} [1+1]$$

$$= \frac{2}{R} \frac{R^{3}}{3}$$

$$y = \frac{Moment of area}{Totalarea}$$

$$=\frac{\frac{2}{3}R^{3}}{\pi R^{2}}$$
$$=\frac{4R}{3\pi}$$

Centroids of different figures

| Shape | Figure | \overline{x} | \overline{y} | Area |
|----------------|---|-------------------|-------------------|---------------------|
| Rectangle | di | $\frac{b}{2}$ | $\frac{d}{2}$ | bd |
| Triangle | + - + - + - + - + - + - + - + - + - + - | 0 | $\frac{h}{3}$ | $\frac{bh}{2}$ |
| Semicircle | A Y Y | 0 | $\frac{4R}{3\pi}$ | $\frac{\pi r^2}{2}$ |
| Quarter circle | y | $\frac{4R}{3\pi}$ | $\frac{4R}{3\pi}$ | $\frac{\pi r^2}{4}$ |

Problem 3: Find the centroid of the T-section as shown in figure from the bottom.



| Area (A _i) | Xi | y _i | A _i x _i | A _i y _i |
|------------------------|----|----------------|-------------------------------|-------------------------------|
| 2000 | 0 | 110 | 10,000 | 22,0000 |
| 2000 | 0 | 50 | 10,000 | 10,0000 |
| 4000 | | | 20,000 | 32,0000 |

$$y = \frac{\sum A_{i}y_{i} A_{1}y_{1} + A_{2}y_{2} 32,0000}{\overline{A_{i1}}}_{2} A + A^{-} 4\overline{0}00$$

Due to symmetry, the centroid lies on Y-axis and it is at distance of 80 mm from the bottom.

Problem 4: Locate the centroid of the I-section.



As the figure is symmetric, centroid lies on y-axis. Therefore, x=0

| Area (A _i) | Xi | yi | A _i x _i | A _i y _i |
|------------------------|----|-----|-------------------------------|-------------------------------|
| 2000 | 0 | 140 | 0 | 280000 |
| 2000 | 0 | 80 | 0 | 160000 |
| 4500 | 0 | 15 | 0 | 67500 |

$$y_{\overline{c}} = \frac{A_{y_i}}{A_{y_i}} = \frac{A_{1}y_1 + A_{2}y_2 + A_{3}y_3}{A_{1} + A_{2} + A_{3} + A_{3} + A_{3}} = 59.71 mm$$

Thus, the centroid is on the symmetric axis at a distance 59.71 mm from the bottom.

Problem 5: Determine the centroid of the composite figure about x-y coordinate. Take x = 40 mm.



 A_1 = Area of rectangle = 12x.14x=168x² A_2 = Area of rectangle to be subtracted = 4x.4x = 16 x² A_3 = Area of semicircle to be subtracted =

$$\frac{\binom{2}{2} \frac{1}{2} \frac{$$

 A_4 = Area of quatercircle to be subtracted =

A₅ = Area of triangle =44
$$\frac{1}{2} \times 6x \rightarrow 4x = 12x^2$$

| Area (A _i) | Xi | yi | A _i x _i | A _i y _i |
|------------------------|--|---|-------------------------------|-------------------------------|
| $A_1 = 268800$ | 7x = 280 | 6x =240 | 75264000 | 64512000 |
| $A_2 = 25600$ | 2x = 80 | 10x=400 | 2048000 | 10240000 |
| $A_3 = 40208$ | 6x =240 | $\frac{4 \times 4x}{} = 67.906$ | 9649920 | 2730364.448 |
| | | 3π | | |
| $A_4 = 20096$ | 10x + 4x - 4x + 10x + 4x + 10x + 4x + 10x + 10 | 8x + 4x - 4x + 4x + 4x + 4x + 4x + 4x + 4 | 9889040.64 | 8281420.926 |
| | (3π) | $\left(3\pi\right)$ | | |
| | = 492.09 | = 412.093 | | |
| $A_5 = 19200$ | $14x + \frac{6x}{1} = 16x$ | $\frac{4x}{2} = 53.33$ | 12288000 | 1023936 |
| | 3 | 3 | | |
| | = 640 | | | |

$$x = \frac{A_{1}x_{1} - A_{2}x_{2} - A_{3}x_{3} - A_{4}x_{4} + A_{5}x_{5}}{A_{1} - A_{1} - A_{1}$$

Problem 6: Determine the centroid of the following figure.



$$A_{1} = \text{Area of triangle} = \frac{1}{2} \times 80 \times 80 = 3200m^{2}$$

$$A_{2} = \text{Area of semicircle} = \frac{\pi d^{2}}{\frac{\pi D^{2}}{2}} = \frac{\pi R^{2}}{2} = 2513.274m$$

$$A_{3} = \text{Area of semicircle} = \frac{\pi d^{2}}{\frac{\pi D^{2}}{2}} = 1256.64m^{2}$$

| Area (A _i) | Xi | yi | A _i x _i | A _i y _i |
|------------------------|----------------|-------------------------|-------------------------------|-------------------------------|
| 3200 | 2×(80/3)=53.33 | 80/3 = 26.67 | 170656 | 85344 |
| 2513.274 | 40 | $-4 \times 40 = -16.97$ | 100530.96 | -42650.259 |
| | | $\overline{3\pi}$ | | |
| 1256.64 | 40 | 0 | 50265.6 | 0 |

$$x = \frac{A_{1}x_{1} + A_{2}x_{2} - A_{3}x_{3}}{A_{1} + A + A_{3}x_{3}} = 49.57mm$$

$$x = \frac{A_{1}x_{1} + A + A_{3}x_{3}}{A_{1} + A + A_{3}x_{3}} = 9.58mm$$

$$y_{c}^{2} = \frac{A_{1}x_{1} + A + A_{3}x_{3}}{A_{1} + A - A_{3}x_{3}} = 9.58mm$$

Problem 7: Determine the centroid of the following figure.



 A_1 = Area of the rectangle A_2 = Area of triangle A_3 = Area of circle

| Area (A _i) | Xi | yi | A _i x _i | A _i y _i |
|------------------------|-----------|----------|-------------------------------|-------------------------------|
| 30,000 | 100 | 75 | 3000000 | 2250000 |
| 3750 | 100+200/3 | 75+150/3 | 625012.5 | 468750 |
| | = 166.67 | =125 | | |
| 7853.98 | 100 | 75 | 785398 | 589048.5 |

$$x_{c}^{=} \sum_{i}^{A_{i}x_{i}} = \frac{A_{1}x_{1} - A_{2}x_{2} - A_{3}x_{3}}{A_{-}A_{-}A_{3}x_{3}} = 86.4mm$$

$$y_{i}^{=} \sum_{i}^{A_{i}y_{i}} = \frac{A_{1}y_{1} - A_{2}y_{2} - A_{3}y_{3}}{A_{-}A_{-}A_{3}x_{3}} = 64.8mm$$

Numerical Problems (Assignment)

1. An isosceles triangle ADE is to cut from a square ABCD of dimension 'a'. Find the altitude 'y' of the triangle so that vertex E will be centroid of remaining shadedarea.



2. Find the centroid of the followingfigure.



3. Locate the centroid C of the shaded area obtained by cutting a semi-circle of diameter 'a' from the quadrant of a circle of radius'a'.



4. Locate the centroid of the compositefigure.



Truss/ Frame: A pin jointed frame is a structure made of slender (cross-sectional dimensions quite small compared to length) members pin connected at ends and capable of taking load at joints.

Such frames are used as roof trusses to support sloping roofs and as bridge trusses to support deck.

Plane frame: A frame in which all members lie in a single plane is called plane frame. They are designed to resist the forces acting in the plane of frame. Roof trusses and bridge trusses are the example of plane frames.

Space frame: If all the members of frame do not lie in a single plane, they are called as space frame. Tripod, transmission towers are the examples of spaceframes.

Perfect frame: A pin jointed frame which has got just sufficient number of members to resist the loads without undergoing appreciable deformation in shape is called a perfect frame. Triangular frame is the simplest perfect frame and it has 03 joints and 03members.

It may be observed that to increase one joint in a perfect frame, two more members are required. Hence, the following expression may be written as the relationship between number of joint j, and the number of members m in a perfect frame.

m = 2j - 3

- (a) When LHS = RHS, Perfectframe.
- (b) When LHS<RHS, Deficientframe.
- (c) When LHS>RHS, Redundantframe.

Assumptions

The following assumptions are made in the analysis of pin jointed trusses:

- 1. The ends of the members are pin jointed(hinged).
- 2. The loads act only at thejoints.
- 3. Self weight of the members isnegligible.

Methods of analysis

1. Method ofjoint 2. Method ofsection

Problems on method of joints

Problem 1: Find the forces in all the members of the truss shown in figure.





 $\tan \theta = 1$ $\Rightarrow \theta = 45^{\Box}$

Joint C

 $S_1 = S_2 \cos 45$ $\Rightarrow S_1 = 40 KN (Compression)$ $S_2 \sin 45 = 40$ $\Rightarrow S_2 = 56.56 KN (Tension)$

Joint D

 $S_3 = 40KN$ (Tension) $S_1 = S_4 = 40KN$ (Compression) Joint

<u>B</u>

Resolving vertically, $\sum_{S_5} V = 0$ $S_5 \sin 45 = S_3 + S_2 \sin 45$







 $\Rightarrow S_5 = 113.137KN$ (Compression)

Resolving horizontally,

 $\sum_{S_6=S_5 \cos 45} H = 0$ $\Rightarrow S_6 = S_5 \cos 45 + S_2 \cos 45$ $\Rightarrow S_6 = 113.137 \cos 45 + 56.56 \cos 45$ $\Rightarrow S_6 = 120KN \text{ (Tension)}$

Problem 2: Determine the forces in all the members of the truss shown in figure and indicate the magnitude and nature of the forces on the diagram of the truss. All inclined members are at 60° to horizontal and length of each member is2m.



Taking moment at point A,

$$\sum M_{A} = 0$$

$$R_{d} \times 4 = 40 \times 1 + 60 \times 2 + 50 \times 3$$

$$\Rightarrow R_{d} = 77.5KN$$

Now resolving all the forces in vertical direction,

$$\sum_{R_a+R_d=40+60+50} V = 0$$
$$\Rightarrow R_a = 72.5KN$$

Joint A

 $\sum_{i=1}^{i} V = 0$ $\Rightarrow R_a = S_1 \sin 60$ $\Rightarrow S_1 = 83.72 KN \text{ (Compression)}$

 $\sum H = 0$ $\Rightarrow S_2 = S_1 \cos 60$



 \Rightarrow S₁=41.86*KN*(Tension)

Joint D

 $\sum_{T} V = 0$ S₇ sin 60 = 77.5 $\Rightarrow S_7 = 89.5KN \text{ (Compression)}$

 $\sum H = 0$ $S_6 = S_7 \cos 60$ $\Rightarrow S_6 = 44.75KN \text{ (Tension)}$

Joint B

 $\sum_{S_1 \le 0} V = 0$ $\Rightarrow S_1 \le 60 = S_3 \ge 60 + 40$ $\Rightarrow S_3 = 37.532KN \text{ (Tension)}$

 $\sum H = 0$ $S_4 = S_1 \cos 60 + S_3 \cos 60$ $\Rightarrow S_4 = 37.532 \cos 60 + 83.72 \cos 60$ $\Rightarrow S_4 = 60.626KN$ (Compression)

Joint C

 $\sum_{S_5} V = 0$ $S_5 \sin 60 + 50 = S_7 \sin 60$ $\Rightarrow S_5 = 31.76KN \text{ (Tension)}$







Plane Truss (Method of Section In case of analysing a plane truss, using method of section after doterming the support reactions a section line is drawn possing through . not more than three members in which forces are unknown, such that the entire rame is cut into two separate parts. Ed. Each part should be in equilibrium under the action of loads, reactions and the forces in the members. Method of section is preferred for the following cases: ers of large truss in which forces in only (i) analy members are required joint fails tostartor proceed with If methodo n grown Dith only two u for not setting a joint analysis forces. Example L. 2010m 10100 10100 IDAN Iden Iter 1 oteri T N B M K 60 T 1 C -7×4 = 28m K Determine the forces in the members it, they, and GI in the trues Ra=Rs= 1 > total downword lood Due to symmetry 2×70: 35KN. To King the section to the left of the cut DEN. D Taking moment about by ZMG = 0. FRHX 951760 +25×12 - I = 10x2+10x6+10x010 67 E C (20+60+100)-420 >> +F# = 25ten, =-69.28 kN. 7 Sin 60

Negative sign indicates that direct opposite 1. e it is compressive in noture Now Resolving all the forces vertically Eyes 10+10+10+ FGH Sin 60 = 35 35-30 TETH = Sin60' FGH = 5:78 KN. / (compressive) Repolving all the force horizontally Ex=0. FFH + fgH cos 60 = fgi FGL = 69.28 + 5.78 cos 60' = 72-17 Knt. ch sion Using method of sections determine the avial forces () in bors 1,2 and 3. F L 52 Teking moment about the joint D. ZMD=0 Th s, xa= Pxh => s1= tension (1) Similarly taking & as the moment contro EME = 0 52×a++×2h=0 -2ph 153 -(-ve sign indicates direction of >> a force Dillbe opposte and it while compressi re in nature Resolving all the force horizontally. Ix=0. coso = - a 1/97 5 52005 x = + V42+12 2) 5, -(Ans)



of (6.3) calculate the relation beth active forces panda for equilibrium of system of bars. The bars are so erranged that they form identical rhambusee. Let 2 = length of each sides & bar. O = angle made by each side of the thambus Distance of from fixed point A: BRCOSD · 22 crso Letthe virtual displacement of P is B-B! B-B'= dry = to (Strand D) and = - 6l sin & dD Similarly the virtual displacement of R is crcl = d12= -28 8100 d0 Applying principled virtual work p. day = R. daz P.(BR 8'n & do) = Q (2R 8'n 8 y + = -3 (Ans) A prismatic bar AB of length l and wt. & stands in a vertical plane > Rb. and is supported by smooth surfaces at B 2/2 Aand B, Usin's principles virtaal work find the magnitude of horizontal + force & applied at A lifthe R baris in equilibrium,

Let she the compressive force in bar CD. consider the part EBDF of the trues under the action of force Rb, Pands Keeping E fixed and giving EB an angular displacement da ZW=0, REXBB' = SXFF! BB'= 2 dd FFI: hdq Roxeda = sxhda B RL. >> S= <u>R50</u> - ci) Now considering whole frame as equilibrium body Rat RS = P. Rbik= P. R = 2 / Rb= 2 00 Substituting the value of Rbin eq. (1) PR 5 = (Ans) ______(3) 44 0.4 (6.15) Using principled virtual work 45 find reactions the for the trues, Let the true is virtual A 145° displaced by an amount dy Ra, ZW=0-Rax AAIS PXDD/ modipada to bigbazar where AAI = DDI = dy her jastrath -> Ra= P mandir righthandhide

02/12/14 0 Momento & Enertico Plane figures The moment of inertia of any plane figure r with respect to x and y ares in its plane are expressed las Las / y 2 dA Lys / nedA Inx and by are also known as seend momento X inertia area about the area as it is distance is squared from corresponding anis. unit Unitof momental inertia of area is copressed as more mont. Momentof Inpertia of Plane figures:i) Restanglo considering a rectangled dy to many man width band depth of, 1 d/2 Momentox inertia about Centroidal acis N-x × - x parollel to the short side d/2 1.e 5 Now considering an elementar e-b]strip of width dy Momento finertra of the elemental strip about centraidal is aris XX Exx = y2 dA = y2 bdy of entire for So momentof inpettia $b \left| \frac{\gamma^{3}}{3} \right|_{-\frac{1}{2}}^{\frac{1}{2}} = b \left[\frac{a^{3}}{2q} + \frac{a^{3}}{2q} \right]_{-\frac{1}{2}}^{\frac{1}{2}}$ J y2 bdy Ín > 1 6d3 >> [xx = d 13 Eyy = 1 mansh at Smilarly

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LAB =
$$\int_{-\frac{1}{2}}^{h} \frac{y^2(h-y)}{h} b dy$$

= $b \left[\frac{y^3}{2} - \frac{y^4}{7h} \right]_{-\frac{1}{2}}^{h} = b \left[\frac{h^3}{2} - \frac{h^4}{7h} \right]$
= $b \left[\frac{h^3}{2} - \frac{h^3}{4} \right] = \frac{bh^3}{12}$
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(2 02/12/14 $=\int_{0}^{\infty}\frac{s^{3}}{2}\left[\theta-\frac{8h_{2}\theta}{2}\right]^{1/2}ds$ $\int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \left(2\pi - \frac{8^{\prime} - 4\pi}{2}\right) dr$ $\left[\frac{84}{8}\right]\left[2\pi-0\right]$ $= \frac{R^{4}}{8} 2\pi = \frac{\pi R^{4}}{4}$ $= \frac{\pi R^{4}}{8} = \frac{\pi R^{4}}{4} = \frac{\pi D^{4}}{64}$ Polar momentox inertia !-Moment of inertia about an aris perpendicular to the plane of area is called potar moment of inertia it may denoted as Torixz 0 LZZE ZordA Radius of Gyrotion !-Radious of system may be defined by a relation K: VI where Karadius of syrotion I = moment of inertia A = cross-sectional area so, we can have the following relations Kax: V Exx kyy = V Lyy LAB

Moment of inputio Theorems of There are two theorems of moment of inertig (a) perpendicular aris theorem parallel and theorem. Perpendicular axis theorem!-Momentox enertia of an area about an adis Ir to Stis plane atany point o is equal to the surs of moments of inertia about any two methody perpendicular acis through the same point o and lying in the plane of area. IXX = IXX + Lyy LXX = Er2dA = Z(227 y2) dA ExedA+EyedA Eze E LXX+ Lyy Parallel axis theorem !-Momentof inertia about an anis TEINIT in the plane of an area is equal to the sum of moment of inertig about a parollel centroidal axis and the productor area and なり square of the distance bet n A the two poralles and. LORB - Love Egg + Ah 2

Moment of inertia of a reatangle continues:
Moment of inertia of a reatangle about 1
it is control all and xx

$$Ixx = \frac{ud^3}{12}$$

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it is (kentrol all and yy 4
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Now moment of inertia of reatangle 4
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poralial and theorem
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Civ) Marcent of inertia of triangle about its bact.
Moment of inertia of triangle about its bact.
Moment of inertia about its centroid

$$\pm A + A^2$$

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 $\pm A + A^2$
 $= \sum \frac{bA3}{12} = \frac{bA3}{12} + \frac{b}{2} bxh x(\frac{b^2}{2})$
 $= \sum \frac{bA3}{12} = \frac{bA3}{12} + \frac{b}{2} bxh x(\frac{b^2}{2})$
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$$\frac{9}{128} = \frac{1}{128} \times \frac{1}{188} \times \frac{1}{987} \times \frac{1}{987} \times \frac{1}{987} \times \frac{1}{987} \times \frac{1}{987} \times \frac{1}{128} \times \frac{1}{1287} \times \frac{1}{128} \times \frac{1}{1287} \times \frac{1}{128} \times \frac{1}{1287} \times \frac{1}{128} \times \frac{1}{128}$$

$$\frac{e_{2}|1_{2}|2_{2}|4}{5}$$
Similarly ME about YY control all ands
$$\frac{1}{1_{Y}} = \frac{5}{1_{2}} \frac{1_{2}}{1_{2}} + 1_{2}50 \times (20.93 - 5)^{2} \frac{3}{2}$$

$$+ \frac{5}{2} \frac{1_{0} \times 7b^{3}}{1_{2}} + 750 \times (47.5 - 20.93)^{2} \frac{3}{2}$$

$$= \frac{10416.66667 + 317206.185}{1_{2}} + \frac{1057562.5 + 529472.655}{1_{2}}$$

$$= \frac{1208658.967 mm^{2}}{1_{2}}$$
Polar momental inortia Izz = Izv + Lyy
$$= \frac{4393217.82}{1_{2}} mm^{2}$$

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$$\frac{1}{1_{1}} centroldal and Yy. An$$$$$$$$$$$$$$$$$$

L

MODULE IV

PARTICLE DYNAMICS AND INTRODUCTION TO KINETICS

- Reallinear Translation :-

In statice, it was considered that the rigid badies are at rest. In dynamice, it is considered that they are in motion, Dynamics is commonly divided into two branches. Kinematics and kinetics,

- In, kinematics we are uncerned with space time relationship of a given motion of abody and not at all with the forces that cause the motion,
- En kinetice we are concerned with finding the kind of motion that as iven body or system of bodies will have under the action of siven forces or with what forces must be applied to produce a desired motion.

Displacement

trem the fig, displacement of a particle x - x can be defined by its x-coordinate, 0 A X measured from the fixed reference point D: - When the particle is to the right of fixed point O, this displacement can be considered poccitive and when it is

towards the sign refet and side it is considered as negative.

General displacement time equation

where fet) = function of times for ecompte [x = c+st]

-

In the above equation C, represents the initial displacement at t = 0, while the constant b shade the rate atwhich displacement increases. It is called uniform rectilinear motion.

second prampipis / x = - 97 where R is propertional to the equareof time. relocity Acceleration The reatilemean motion of a particle is defined Eram plp by the displacement - time equation x = ko- lot + bat? construct displacement - time and velocity diagram for this motion and find the displacement (and velocity attime te = 25. No = 750 mm, to = 500 mm/s (a = 0:125 m/s2 motron i's The equation of 20-00++ 20+2 - c) 7 : v= dx = -votat substiting no, no and ain equation (1) 2 - 75-500-Difflacemond velocity time. Fime

-1 A beellot leaveethe muxile of o sun with relocity 10 = 750 m/s. Accuming constant acceleration (from breech to muxxie find time t occurpred by the bullet in travelling through gun barrel which is 750 mon long. initial velocity of bullet un o final relative of bollet N=750 m/1, total distance sa 0.75 m. +:2 Wehave V2-42: 200, => v2 = 2as => a = v2 = 7502 270.74 = 375000 / Lee 2 V= letat Again + 750 = 375000 x + >7 t= 750 c 0.002 see. Astone is dropped into well and falls vertically with constant acceleration g= 9. sportsee (The sound of impact of stone in the bottom of woll is heared after 6.5 see. if relating of sound is 336 m/r. Knew deep is the over ?. V= 336m/see 1.2+5: depth of well the time taken by thestone into the well to time taken by the sound to be heared. total time t = (fipt2) = Bib See, Now s= let p is ste >> SIO + 1 st2 >> += /25 When the sound travels with uniform relocity GE Vt2 or t2 = 1

25 + 5 = . B'5 $\frac{25}{25} = \left(6.5 - \frac{5}{336} \right)$ 336 9.87 (25 -2189 5 0.0291 (2184-5) = 0.0291 (4769856 + s2 - 4368 S 138802.809 + 0.029, 12 - 1257. 10885 0:029152-129.10868+138802, sog =0 >> 5 = 0.20385 = 42.25 + 0.0000 88552 - 0.03865 5-2 174 0.0000 885552 - 0.16586 + 42.2520 52 17. 31m, Arope ABis attached at B to a small blockox A2 negligible dimpositors and possessiver a prelleg C sothat it's free end A hanks ison above Browned when the block reets on the floor. The end A of the rope is moved horizontolly in act. link by a man walking with a uniform velocity to plot the velocity time drag ram - 3m/L. (b) find the time trequired for the block to reach the polley if h = 4.5m, pully dimension are negligible, Aparticle starts firm nestand movee along q A3 stilling with constant acceleration a. Efit acquires a velocity u= 3 m/s. after having travelled a distance so 7.5 m. find magnitude of acceleration.

20/11/2014 Principles of Dynamics Newton's law of motion! first law! Everybody continues in it's state of restor of ceniform motion in actraight line except in so for as it may be compelled by force to change that state, Seemd Laed : + The acceleration of a given particle is propertional to the force opplied to it and take place in the direction of the straight line in which the force oets. Third law To every action there is always on equal and contrary reaction or the meteral actions of any two bodies are lationage equal and oppositely directed General Equation of Motion of a Particla! rona = f Dioferential equation of Reatilinear motion :. Differential form of equation for rectilinear motion can be epressed as W x = x x'= acceleration Where X = Receltant acting force. zxamplp For the engine shown in B Marcine fig, the established It. of piston and priston rod W= 450N, cronk rodius r = 250mm and renitor m speed of rotation n= 120 ppm, potermine the magnitude of resultant force acting in priston ca) at afterme position and at the middle position
priston has a simple harmonic motion represented
displacement-time equation

$$x = r \cos 3t$$
 — cr)
 $w = \frac{3\pi n}{60} = \frac{3\pi n \sin 20}{60} = 4\pi rad/s$.
 $\dot{w} = -rw^2 \cos 3t - c^2$
Differential equation of notion
 $\frac{1}{10} \frac{1}{12} = x$
 $\frac{1}{12} \frac{1}{12} \frac{$

Wa = (W-+) (w-2)a = - (w-2) Wa + (W- R) 9 2 W- p8 + p8 - (RI-R) 2 Watula-Ra e R 27 $\frac{2}{\sqrt{\alpha}} = \frac{2}{\sqrt{\alpha}} \frac{2}{\sqrt{\alpha}} \frac{2}{\sqrt{\alpha}} \frac{1}{\sqrt{\alpha}}$ A wir W = 9450N is supported in a vertical plane 1-1 by string and pulleys arranged showing fig. If the free and hoy at the string is pulled vertically downword with constant accoloration as 18m/s2 find kension s in the string. Differential squation of motion for the system is $2s - W = \frac{W}{g} \times \frac{a}{2}$ W + Wa 28 => 25= 2 27 29 $W(1+\frac{q}{2g})$ W (1+ 9 2 (1+ 28) +450 (1+ 18 2×9.81 4266 . 28 N. 2

$$\frac{W_{R}}{S} = \frac{W_{R}}{S} + \frac{W_{R}}{S} = \frac{W_{R}}{S} + \frac{W_{R}}{S} - \frac{W_{R}}{S} + \frac{W_{R}}{S} = \frac{W_{R}}{S} + \frac{W_{R}}{S} - \frac{W_{R}}{S} + \frac{W_{R}}{S} = \frac{W_{R}}{S} + \frac{W_{R}}{S} +$$

An elevator of gross with = 4450N starts to move. upideral direction with a constant acceleration and acquires avelocity & : 15m/s; after travelling a distance = 1.80m, tind tensile force sin the Cable during it's motion. -V: 15m/s. メミレシカの W= 4450N. V: 18m/s. initial velocity u: 0 glistance tratelled x=1.5m, 1N=4450 N. $S-W = \frac{W}{R} \cdot q$ $\gamma s = w + \frac{w}{s} a = w \left(1 + \frac{a}{s} \right)$ Now applying equation of bing to atre N2-42= 2as 27 182-0 = 2a×118 27 a 2 182 5 90 m/22 substituting the value of a 4150 (17 90) = 45275.7 S 2 A-7) A train Deighing 1870H without the locomotive starts to move with constant acceleration along q straight track and in first 600 acquires a velocity of 56 Kmph, Determine the tensions in draw bar beth locomotive and train if the air resistance is 0.005 times the oft. of the train. V: 56 Kmph = 15.56 m/2. - a > H=0 -> S F= D.DOSW < W=1870N.

S-F= Wig 27 5= 0:005W+ Wa from eq. of elners attice. v: utat => a = (15:56-0) = 0:26 m/see 2 substituting the value of a in eq. (1) W (0.005 + 9 g 5 = 1870 (0:005 + <u>6:26</u>) = [5 8.9 KN. 1 A with up is attached to the and of actuall flexible rope of dia, d= 6:25mm, and is raised vertically by Dinding the rope on a real. if the real is turked uniformity at a rate of 2 rpc. what will be the tension (in rope . dia of rope d= 6:25mm = 0:00825m, No of revolutions N= 2 rps. let x = initial rodius of real. t: there taken for M revolutions. -0 Metrodius after + see. 2 R= [x+(N+d)] Now maan velocity N= Rw N= 2MN. · V= (x+ N+d) 2771 acceleration of sope & a = di $a = \frac{d}{dt} \left[2\pi n x + a\pi n^2 + d \right] = a\pi n^2 d$ S-W= W, q =>S= W+ Wq = W(1+ q) => S= W (1+ 29TN2 +) =+

$$\frac{1}{2} = W \left(1 + \frac{2\pi \sqrt{2^2 \times 0.6\pi 22^{N}}}{9.57} \right)$$

$$=$$

$$\frac{1}{2}$$

$$\frac{1$$

25/11/14 112 Princip Differential equation of motion (rectilinear) can be written as X-mx=D - (1) Where x = Resultant of all applied force in the direction of rotion m: mass of the particle The above equation may be treated as equation of dynamic equilibrium. To supress this equation, in addition to the real force acting on the porticle a fictitious force mix is required to be considered. This force is equal to the productor make of the particle and it is acceleration and directed opposit direction, and is called the inertia force of the particle. - このに = - デ えの $= -\frac{W}{0}\ddot{z}$ W - total weight of the body Where so the equation of dynamic equilibrium can be expressed as! $\Sigma x_i + \left(-\frac{W}{g}\ddot{z}\right) = 0$ Example for the Reample shown considering the motion of pellagy as shown by the arrow mork 711911 wehave upwhed acceleration to for W2 and downward acceleration is for W, - corresponding inertia force and their direction are indicated by dotted TI line. +w, - By adding inertra force to the real forces (such as W, W2 and tension in strings) we obtain, for each particle, a system of 100,76 man forces in equilibrium. The equilibrium equation for the entire eyelem with not S

 $W_2 + m_2 \ddot{x} = W_1 - m_1 \dot{z}$ $= (m_1 + m_2) \ddot{x} = (W_1 - W_2) = / \ddot{z} = \frac{W_1 - W_2}{(W_1 - W_2)}$

body is moving in upperd direction a ropo. so the equation of dynamic equilibrium considering the real and inertia forces. S-W-Ba=0, so tensile force in rope W 5 = W (1+ a R 少业 光 Find tensions in the string during motion of the cyclem 0 Coo) if W; = 900 N, W2= 450 N. The pebet the Inclined plane and block W1 = 0.2 When W, moves doonward in the inclined plane with an ee acceleration a, then acceleration of H2 = Considering dynamic equilibrium of Mi, from DI Alembertis principle (W, Sin 45'- 12N1 - 5) - W1 a = 0 W1 a = W, Sin 45' - Jen - S Wisin45 - MW, 10545-5 $= \left(900 \times \frac{1}{V_2} - 0.2 \times 900 \times \frac{1}{V_2} - 5 \right) \frac{9.81}{9.2}$ a => 900 = (636.4 - 127.28 - 5) 0.0109=> $a = \frac{693676}{693676} - 1.967352 + 0.0109 - 0.0109 - 0.0109$ Similarly for for warght W2 $2s - W_2 - \frac{W_2}{R} \frac{a}{2} = 0$ $W_2\left(1+\frac{q}{2g}\right)$ 1= 25 $\frac{450}{2}\left(1+\frac{9}{19.62}\right)$ substituting the value of sin eq. ci

2571/14 693676-1.387352-0.0109 225+11.969 a : = 6-93 5.549408 - 2.4525 - 0.1249149 3.096908 - 0-1299149 => a: 2.75 m/s2 Two weights P and & are connected by the arrangement 0.2 shadn in fig. Neglecting friction and inertia of pudley and cord find the acceleration a of wit- & Assume 7= 175 N, 8= 133.5 N. Applying D' Alembert 15 principle for Q 111414 R-5- Ra=0 = $y = a(1 - \frac{a}{3}) - c1)$ = 133.5(1- q.) > yine b'Alembertis principle to P Applying P=25= $2S - P - \frac{Pa}{2g} = 0$ => 28 = \$ (1+ 9 = 28 => == == (1+ a) C2 178 (17 9 / - $\left(1-\frac{a}{9.51}\right) = 89\left(1+\frac{a}{19.62}\right)$ 133.5 89 + 4.5369 133.5- 13:6089 >/a = 2.95 m/s2 CAns Accuming the car in the fight to have a velocity Emps findehortest distance in which it's with constant decelocation without disturbing the block. potation = orbon, he aig m 31 M= 0.5



$$\frac{e_{\mu} \cdot e_{\mu} \cdot e_$$

27/11/2014 (1) Momentum and impulse We have the differential squation of rectilinear motion of a particle W x = X Above equation may be written as W di = X d(Wx)= Xd+ Dr -CI) En the above equation we dill alcome force x as a function of time represented by a force time diagram. The righthand side of egice). is then represented by the area of shaded elemental strip of ht X and Didte dt. This quantity i.e (Xd+) is called imposed of the force Adre X in time dt. The expression on the left hand side (wind) is called momentum of of the sepression porticle, sothe eq. (1) represents the differential change in momentum of a toarticle in time dt. Lategrating equal) we have $W = x + C = \int^{t} x dt | - c2)$ where C is a constant of integration Now assuming an initial moment, 420, the particle has an initial velocity to C= - 1/20 - (3 50 So equestion (2) beeringes Itxdt

from equation (2) it's clear that the total cha momentum of a particle during a finite interval of the is equal to the impulse of acting force, in other words Q Company (f.dt=d(mv) where mx v= momentum Regoidastran 0-1 A man of wit 712 M stands in a boat so that he is 4.5 m from a pier on the shore. He walks 2.4 m in the boat towards the pier and then stops. How for from the pier with he be at the end of time. Wt of boat is vor g 890m. whoyman inl, = 712 M /sV wt of boat Wa = squap Let vo is the initial verseity of man and fistime then Not = x >> vote airin -> vo = (2.4) m/s. let V = velocity of boat towards right according to concervation of momentum W, Vo = (W, +W2) V => V= (W, TW2) distance covared by boat $s : V : t = \frac{W, V_0}{(W, T W_2)} : t$ 712× a-4 . + = 1.067 m => 5 = × (712+890)

od sidek wit 22.25 M rostsona 50 Awo revolver bullet weighing oiled is shot surface. A into the side of block . If the block Lorizontelly ZZIR attains & relocity of 3m/3 whatis velocit WI. of wood block M, = 22.25 Nd. Wtion wollet Was out of. V= Bron /s. velocity of velocity Mouzzle conservation of momentum According to M, K= 1424 = (11, + 12) (22.25+0.14 479.98 m/s. Conservation of momentum impulses due to When the the system r R M John conserved the momentum o Z St X at = 0 When Z $\frac{W}{s}$) $x_{a} = Z\left(\frac{W}{s}\right)\dot{x}_{i}$ initial momentum tinal momentum =

Cervilinear Translation When is moving porticle describes a worved poth it is said to have curvilinear motion. Displacement consider a particle Pinaplaneona Lerred poth. Ax Todefine the particle we need two coordinate randy as the particle mones, 0 there evoratino te marce change with time and the displacement time equations are x= f(+) y= f2 (+) The motion of porticle con also be coproceed as 5 = f, (+) 7= f (* represents the equation of path 04 7=F(x) where and s=fi(t) sives displacements measured along the path as a function of time. velocity :-Considering an infiniteermal time difference from t to + + 1+ during which the porticle move from ptop along it's path. velocity of porticle may be expressed as then AE At Var = 41 aregage velocity along Dav /x 1 38 continates rand y Var), =

It can also be pepressed, al Oz = dr = x oy - dy = y cothe total velocity may be represented and $\cos(0, x) = \frac{x}{12}$ and $\cos(0, y) = \frac{y}{12}$ where \$ (0,x) and (0,y) denotes the ons bet the direction of velocity vector to and the coordinate ane Acceloration :-The eccleration porticles may be described as ax = dr = r ay = dy = y L'é is also known as instantaneous acceleration Total acceleration a: / 2+ y2 Considering particular path for above care. x: ressedt y=rsignedt. 2742=12 2= - rwinnot y= rwcosit $\phi = \sqrt{i^2 + j^2}$ 2= -rw2 wit y= -rw2 sinwf a = /x = + y 2

D'Alemberts principle in curvilinear Motion



Condition for skidding !-Let W = wt of vehicle R, R2 = . realtions at wheel F = frictional force. W. UP = inpritia force skidding takes place when the firstional forces reaches limiting value i.e F= pent permissible speed to avoid skid-ling Thenmoum gr B D= The distance betn inner and outer wheel is equal to the gauge of railway track and supressed as by. er G 50 Designed speed and ansient Broking Z of all the forces in the Inclined plane W u2 coso - W Sind =0 W u2 122 => tand = gr R2 RI d Relation befor the angle of broking and designed speed 122 tondo 13 Q8

$$\frac{1}{12} = \frac{1}{12} + \frac{1}{12}$$

D' Alembert's Principle in Curvilinear Motion Equation of motion of a porticle maybe written as X-mi=0 -- c1) Y - ~ ÿ=0 find the proper super elevation 'e' for 07.2 m 0.2 high day curve of radius r= boom in order that a car travelling with aspeed of 80 Kmph will have no tendency toskid side dise. R TX b=7.2m r= 600m V= 80Kmph= 22.23 m/2. Resolving along the inclined plane ' WEIND = W. V2 US of >> tond = - v2 rg from the permetry sind = e, since d is norgenall let sind a fag $\frac{V^2}{r_g} = \frac{g}{b} = \frac{5v^2}{r_g} = \frac{7 \cdot 2 \times 22 \cdot 23^2}{\frac{5v^2}{r_g}} = \frac{7 \cdot 2 \times 22 \cdot 23^2}{\frac{5v^2}{r_g}}$ 2 0.60 qm (Ans)

)4

acing car travels around a circu soom radius with aspeed of 884 kmph. 01 what and a shreld the floor of the track make with horizontal in order to safeguard against skidding. velocity &: 384 kersph 2 = 300m = 106.67 m/s. we have angle of breking tand: us 106.672 of de tant 2 75.50 (supp) Two bolls of wit Wa = 44.57 and Ws = 66.751 are connected by an elastic string and supported on a timeto le as shown. When the furnite we is of rat, the tension in the string is s= 222.5N and the balls event this same force on each of the stops A and B. What forces will they piert on the stops when the tiern table is rotating uniforming assoct the vertical aurs CD at 60 spm 2 250 mont 250 mm Wehave; Wo = 44.50 WS = 66.751 5 = 222.5N m= 60 spm, radices of rotation r, r2=0:25m Now angular Walved teg ent : an ADO & 211 red C 02



Let
$$d = ansular acceloration of the accomply (3)
L = more moment of inartia of the accomply
I = Let Md^2 (transfer formula)
 $I = Let Md^2$ (transfer formula)
 $I = Let Md^2$ (transfer formula)
 $I = bridden + A = -\frac{1}{2} \times \frac{2\pi}{7.51} \times 1^2 + \frac{2\pi}{9.51} \times (0.5)^2$
 $= 6.7968$
mass ML of cylinder about A
 $2 -\frac{1}{2} -\frac{5\pi}{9.51} \times 0.2^2 + \frac{5\pi}{9.51} \times 1.2^2$
 $= 74.4$
 $ME of the cystem = 6.7968 + 74.4 = 51.2097$
Rotational moment as boset A
 $M_{\pm} = 200\times0.5 + 500\times1.2 = 700 + 1 \text{ m}$,
 $M_{\pm} = \frac{760}{81,2097} = \frac{6.6197}{100} \log \log 48 \text{ is}$
 $vertical and $= 1, d = 0.5 \times 8.6197$
 $= 4.31 \text{ m/s}$.
Similarly instantaneous acceleration of rod AB is
 $vertical and $= 1.2\times 8.697$
 $= 10.29 \text{ m/s}$.
Applying Diflembort's dynamic equilibrium
 $R_{A} = 200+500 - \frac{200}{7.91} \times 4.31 - \frac{570}{9.51} \times 10.34$$$$$

MODULE V MECHANICAL VIBRATIONS

Definitions and Concepts

Amplitude :Maximum displacement from equilibrium position; the distance from the midpoint of a wave to its crest or trough.

Equilibrium position: The position about which an object in harmonic motion oscillates; the center of vibration.

Frequency: The number of vibrations per unit of time.

Hooke's law: Law that states that the restoring force applied by a spring is proportional to the displacement of the spring and opposite in direction.

Ideal spring: Any spring that obeys Hooke's law and does not dissipate energy within the spring.

Mechanical resonance: Condition in which natural oscillation frequency equals frequency of a driving force.

Period: The time for one complete cycle of oscillation.

Periodic motion: Motion that repeats itself at regular intervals of time.

Restoring force: The force acting on an oscillating object which is proportional to the displacement and always points toward the equilibrium position.

Simple harmonic motion: Regular, repeated, friction-free motion in which the restoring force has the mathematical form F = -kx.

Simple Harmonic Motion

A pendulum, a mass on a spring, and many other kinds of oscillators exhibit a special kind of oscillatory motion called Simple Harmonic Motion (SHM).

SHM occurs whenever :

i.

| | | | о <u>ни</u> т |
|--------------------|--------------------|---------------------|----------------------------------|
| ere is a restoring | torce proportional | to the displacement | from equilibrium: $F \propto -x$ |
| 0 | , I I | L | 1 |

ii.

he potential energy is proportional to the square of the displacement: $PE \varpropto x^2$

iii.

he period T or frequency f = 1 / T is <u>independent</u> of the <u>amplitud</u>e of the motion.

iv.

he position x, the velocity v, and the acceleration a are all sinusoidal in time.

h

t

t

t



(*Sinusoidal* means sine, cosine, or anything in between.) As we will see, any one of these four properties guarantees the other three. If one of these 4 things is true, then the oscillator is a simple harmonic oscillator and all 4 things must be true.

Not every kind of oscillation is SHM. For instance, a perfectly elastic ball bouncing up and down on a floor: the ball's position (height) is oscillating up and down, but none of the 4 conditions above is satisfied, so this is not an example of SHM.

A mass on a spring is the simplest kind of Simple Harmonic Oscillator.



Notice that Hooke's Law (F = -kx) is condition i : restoring force proportional to the displacement from equilibrium. We showed previously (Work and Energy Chapter) that for a spring obeying Hooke's Law, the potential energy is U = $(1/2)kx^2$, which is condition ii. Also, in the chapter on Conservation of Energy, we showed that F = -dU/dx, from which it follows that condition ii implies condition i. Thus, Hooke's Law and quadratic PE (U $\propto x^2$) are equivalent.

We now show that Hooke's Law guarantees conditions iii (period independent of amplitude) and iv (sinusoidal motion).

We begin by deriving the *differential equation* for SHM. A differential equation is simply an equation containing a derivative. Since the motion is 1D, we can drop the vector arrows and use sign to indicate direction.

$$F_{net} = ma$$
 and $F_{net} = -kx \implies ma = -kx$
 $a = dv/dt = d^2x/dt^2 \implies \frac{d^2x}{dt^2} = -\frac{k}{m}x$

The constants k and m and both positive, so the k/m is always positive, always. For notational convenience, we write $k / m = \omega^2$. (The square on the ω reminds us that ω^2 is always positive.) The differential equation becomes

$$\frac{\mathrm{d}^2 x}{\mathrm{d} t^2} = -\omega^2 x$$

(equation of SHM)

This is the *differential equation* for SHM. We seek a solution x = x(t) to this equation, a function x = x(t) whose second time derivative is the function x(t) multiplied by a negative constant $(-\omega^2 = -k/m)$. The way you solve differential equations is the same way you solve integrals: you *guess* the solution and then check that the solution works.

Based on observation, sinusoidal solution: $x(t) = A \cos(\omega t + \varphi)$,

where A, ϕ are <u>any</u> constants and (as we'll show) $\omega = \sqrt{\frac{k}{m}}$.

A = amplitude: x oscillates between +A and -A

 φ = phase constant (more on this later)

Danger: ωt and φ have units of radians (not degrees). So set your calculators to radians when using this formula.

Just as with circular motion, the angular frequency ω for SHM is related to the period by

$$\omega = 2\pi f = \frac{2\pi}{T}, T = \text{period.}$$

(What does SHM have to do with circular motion? We'll see later.)

Let's check that $x(t) = A \cos(\omega t + \varphi)$ is a solution of the SHM equation.

| Taking the first derivative dx/dt, | we get $v(t) = \frac{dx}{dt} = \frac{dx}{dt}$ | $-A \omega \sin(\omega t + \phi)$ |) . |
|------------------------------------|--|---|---------------------------------|
| | $d \frac{dt}{\cos(\omega t + \omega)} =$ | $d\cos(\theta) d\theta$ |). |
| Here, we've used the Chain Rule: | $\frac{1}{dt}$ | $\frac{d\theta}{d\theta}\frac{dt}{dt},$ | $(\theta = \omega t + \varphi)$ |
| | $= -\sin\theta \cdot \omega = -\omega\sin(\omega t + \varphi)$ | | |

Taking a second derivative, we get

$$a(t) = \frac{d^{2}x}{dt^{2}} = \frac{dv}{dt} = \frac{d}{dt} \left(-A \omega \sin(\omega t + \varphi)\right) = -A \omega^{2} \cos(\omega t + \varphi)$$
$$\frac{d^{2}x}{dt^{2}} = -\omega^{2} \left[A \cos(\omega t + \varphi)\right]$$
$$\frac{d^{2}x}{dt^{2}} = -\omega^{2} x$$

This is the SHM equation, with $\omega^2 = \frac{k}{m}$, $\omega = \sqrt{\frac{k}{m}}$

We have shown that our assumed solution is indeed a solution of the SHM equation. (I leave to the mathematicians to show that this solution is unique. Physicists seldom worry about that kind of thing, since we know that nature usually provides only one solution for physical systems, such as masses on springs.)

We have also shown condition iv: x, v, and a are all sinusoidal functions of time:

 $\mathbf{x}(t) = \mathbf{A}\cos(\omega t + \varphi)$

- $v(t) = -A\omega \sin(\omega t + \phi)$
- $a(t) = -A \omega^2 \cos(\omega t + \varphi)$

The period T is given by $\omega = \sqrt{\frac{k}{m}} = \frac{2\pi}{T} \Rightarrow T = 2\pi\sqrt{\frac{m}{k}}$. We see that T does not depend on the amplitude A (condition iii).

Let's first try to make sense of $\omega = \sqrt{k/m}$: big ω means small T which means rapid oscillations. According to the formula, we get a big ω when k is big and m is small. This makes sense: a big k (stiff spring) and a small mass m will indeed produce very rapid oscillations and a big ω .

A closer look $atx(t) = A cos(\omega t + \varphi)$

Let's review the sine and cosine functions and their relation to the *unit* circle. We often define the sine and cosine functions this way:



This way of defining sine and cosine is correct but incomplete. It is hard to see from this definition how to get the sine or cosine of an angle greater than 90°.

A more complete way of defining sine and cosine, a way that gives the value of the sine and cosine for *any* angle, is this: Draw a *unit* circle (a circle of radius r = 1) centered on the origin of the x-y axes as shown:

Define sine and cosine as

 $\cos \theta = \frac{adj}{hyp} = \frac{x}{1} = x$ $\sin \theta = \frac{opp}{hyp} = \frac{y}{1} = y$

This way of defining sin and cos allows us to compute the sin or cos of *any* angle at all.

For instance, suppose the angle is $\theta = 210^{\circ}$. like this:

The point on the unit circle is in the third and y are negative. So both $\cos\theta = x$ and te the sin or cos of Then the diagram looks quadrant, where both x $\sin\theta = y$ are negative (more than once around

For any angle θ , even angles bigger than 360°

the circle), we can always compute sin and cos. When we plot sin and cos vs angle θ , we get functions that oscillate between +1 and -1 like so:

point (x, y)



We will almost always measure angle θ in radians. Once around the circle is 2π radians, so sine and cosine functions are periodic and repeat every time θ increases by 2π rad. The sine and cosine functions have exactly the same shape, except that sin is shifted to the right compared to $\cos by \Delta \theta = \pi/2$. Both these functions are called *sinusoidal* functions.



The function $\cos(\theta + \phi)$ can be made to be anything in between $\cos(\theta)$ and $\sin(\theta)$ by adjusting the size of the *phase* ϕ between 0 and -2π .

$$\cos \theta$$
, $(\phi = 0) \rightarrow \sin \theta = \cos \left(\frac{\theta - \pi}{2} \right)$, $(\phi = -\pi/2)$

The function $cos(\omega t + \phi)$ oscillates between +1 and -1, so the function $Acos(\omega t + \phi)$ oscillates between +A and -A.



Why $\omega = \frac{2\pi}{T}$? The function $f(\theta) = \cos\theta$ is periodic with period $\Delta \theta = 2\pi$. Since $\theta = \omega t + \varphi$, and φ is some

constant, we have $\Delta \theta = \omega \Delta t$. One complete gycle of the cosine function corresponds to $\Delta \theta = 2\pi$ and $\Delta t = T$, (T is the period). So we have $2\pi = \omega T$ or $\omega = \frac{T}{T}$. Here is another way to see it: $\cos(\omega t) = \cos \left| 2\pi \frac{T}{T} \right|$ is periodic

with period $\Delta t = T$. To see this, notice that when t increases by T, the fraction t/T increases by 1 and the fraction $2\pi t/T$ increases by 2π .



Now back to simple harmonic motion. Instead of a circle of radius 1, we have a circle of radius A (where A is the amplitude of the Simple Harmonic Motion).

SHM and Conservation of Energy:

Recall $PE_{elastic} = (1/2) k x^2 =$ work done to compress or stretch a spring by distance x.

If there is no friction, then the total energy $E_{tot} = KE + PE = constant$ during oscillation. The value of E_{tot} depends on initial conditions – where the mass is and how fast it is moving initially. But once the mass is set in motion, E_{tot} stays constant (assuming no dissipation.)

At any position x, speed v is such that $\frac{1}{2}mv^2 + \frac{1}{2}kx^2 = E_{tot}$.

When $|\mathbf{x}| = \mathbf{A}$, then $\mathbf{v} = \mathbf{0}$, and all the energy is PE: $\mathbf{E}_{\mathbf{w}} + \mathbf{E}_{(1/2)\mathbf{k}\mathbf{A}^2} = \mathbf{E}_{tot}$ So total energy $\mathbf{E}_{tot} = \frac{1}{2}\mathbf{k}\mathbf{A}^2$

When x = 0, $v = v_{max}$, and all the energy is KE: $\mathbf{E}_{(1/2)\mathbf{m}v_{max}^2} + \mathbf{P}_0 = E_{tot}$ So, total energy $\mathbf{E}_{tot} = \frac{1}{2}\mathbf{m}\mathbf{v}_{tot}^2$





So, we can relate v_{max} to amplitude A : $PE_{max} = KE_{max} = E_{tot} \Rightarrow \frac{1}{2} k A^2 = \frac{1}{2} m v_{max}^2 \Rightarrow$

$$v_{max} = \sqrt{\frac{k}{m}} A$$

Example Problem: A mass m on a spring with spring constant k is oscillating with amplitude A. Derive a general formula for the speed v of the mass when its position is x.

Answer:
$$v(x) = A \sqrt{\frac{k}{m}} \sqrt{1 - \left(\frac{x}{A}\right)^2}$$

Be sure you understand these things:


Pendulum Motion

A simple pendulum consists of a small mass m suspended at the end of a massless string of length L. A pendulum executes SHM, <u>if</u>the amplitude is not too large.



The restoring force is the component of the force along the direction of motion:

 $F_{restore} = -\begin{pmatrix} mg \\ -L \end{pmatrix} x$ is exactly like Hooke's Law $F_{restore} = -k x$, except we have replaced the constant k with

another constant (mg / L). The math is exactly the same as with a mass on a spring; all results are the same, except we replace k with (mg/L).

$$T_{spring} = 2\pi \sqrt{\frac{m}{k}} \implies T_{pend} = 2\pi \sqrt{\frac{m}{(mg/L)}} = 2\pi \sqrt{\frac{L}{g}}$$

Notice that the period is independent of the amplitude; the period depends only on length L and acceleration of gravity. (But this is true only if θ is not too large.)

SHM and circular motion

There is an exact analogy between SHM and *circular motion*. Consider a particle moving with constant speed v around the rim of a circle of radius A.

The x-component of the position of the particle has *exactly* the same mathematical form as the motion of a mass on a spring executing SHM with amplitude A.



Angular velocity $\mathbf{x} = \frac{d\theta}{dt} = \text{const}$ $\theta = \omega t \text{ so}$

This same formula also describes the *sinusoidal* motion of a mass on a spring.

That the same formula applies for two different situations (mass on a spring & circular motion) is no accident. The two situations have the same solution because they both obey the same equation. As Feynman said, "The same equations have the same solutions". The equation of SHM is $\frac{d^2x}{dt^2} = -\omega^2 x$. We now show that a

particle in circular motion obeys this same SHM equation.

Recall that for circular motion with angular speed ω , the acceleration of a the particle is toward the center and has

magnitude $|\ddot{a}| = \frac{V^2}{R}$. Since $v = \omega R$, we can rewrite this as $|\ddot{a}| = \frac{(\omega R)^2}{R} = \omega^2 R$



Example A mass of 0.5 kg oscillates on the end of a spring on a horizontal surface with negligible friction according to the equation $x = A\cos(\omega t)$. The graph of *F vs. x* for this motion is shown below.



The last data point corresponds to the maximum displacement of the mass. Determine the

(a) angular frequency ω of the oscillation,

(b) frequency *f* of oscillation,

(c) amplitude of oscillation,

(d) displacement from equilibrium position (x = 0) at a time of 2 s.

Solution:

(a) We know that the spring constant k = 50 N/m from when we looked at this graph earlier. So,

$$\omega = \sqrt{\frac{k}{m}} = \sqrt{\frac{50 N/m}{0.5 kg}} = 10 \frac{rad}{s}$$

(b) $f = \frac{\omega}{2\pi} = \frac{10 rad/s}{2\pi} = 1.6 Hz$

(c) The amplitude corresponds to the last displacement on the graph, A = 1.2 m.

(d) $x = A\cos(\omega t) = (1.2 m)\cos[(10 rad / s)(2s)] = 0.5 m$

Example

A spring of constant k = 100 N/m hangs at its natural length from a fixed stand. A mass of 3 kg is hung on the end of the spring, and slowly let down until the spring and mass hang at their new equilibrium position.



(a) Find the value of the quantity *x* in the figure above. The spring is now pulled down an additional distance *x* and released from rest.

(b) What is the potential energy in the spring at this distance?

(c) What is the speed of the mass as it passes the equilibrium position?

- (d) How high above the point of release will the mass rise?
- (e) What is the period of oscillation for the mass?

Solution:

(a) As it hangs in equilibrium, the upward spring force must be equal and opposite to the downward weight of the block. F_s

$$F_{s} = mg$$

$$kx = mg$$

$$x = \frac{mg}{k} = \frac{(3kg)(10m/s^{2})}{100 N/m} = 0.3$$

mg

(b) The potential energy in the spring is related to the displacement from equilibrium position by the equation $U = \frac{1}{2}kx^2 = \frac{1}{2}(100 N/m)(0.3m)^2 = 4.5 J$

(c) Since energy is conserved during the oscillation of the mass, the kinetic energy of the mass as it passes through the equilibrium position is equal to the potential energy at the amplitude. Thus,

$$K = U = \frac{1}{2}mv^{2}$$
$$v = \sqrt{\frac{2U}{m}} = \sqrt{\frac{2(4.5J)}{3kg}} = 1.7 \, m/s$$

(d) Since the amplitude of the oscillation is 0.3 m, it will rise to 0.3 m above the equilibrium position.

(e)
$$T = 2\pi \sqrt{\frac{m}{k}} = 2\pi \sqrt{\frac{3 \, kg}{100 \, N \, / m}} = 1.1s$$

Example

A pendulum of mass 0.4 kg and length 0.6 m is pulled back and released from and angle of 10° to the vertical.

(a) What is the potential energy of the mass at the instant it is released. Choose potential energy to be zero at the bottom of the swing.

(b) What is the speed of the mass as it passes its lowest point?

This same pendulum is taken to another planet where its period is 1.0 second. (c) What is the acceleration due to gravity on this planet?

Solution

(a) First we must find the height above the lowest point in the swing at the instant the pendulum is released.

Recall from chapter 1 of this study guide
that
$$h = L - L \cos\theta$$
.
Then
 $U = mg(L - L \cos\theta)$
 $U = (0.4kg)(10m/s^2)(0.6m - 0.6m\cos 10^\circ) = 0.\frac{h}{10^\circ}$

(b) Conservation of energy:

$$U_{\text{max}} = K_{\text{max}} = \frac{1}{2}mv^{2}$$

$$v = \sqrt{\frac{2U}{m}} = \sqrt{\frac{2(0.4J)}{0.4 \, kg}} = 1.4 \, m/s$$

$$(c) \begin{cases} T = 2\pi \sqrt{\frac{1}{g}} \\ g = \frac{4\pi^{2}L}{T^{2}} = \frac{4\pi^{2}(0.6m)}{(1.0 \, s)^{2}} = 23.7 \frac{m}{s^{2}} \end{cases}$$

COMPOUND PENDULUM

AIM:

The aim of this experiment is to measure g using a compound pendulum.

YOU WILL NEED:

WHAT TO DO:

First put the knife edge through the hole in the metre rule nearest end A, and record the time for 10 oscillations. Hence work out the time for one oscillation (T).

Repeat this for each hole in the ruler for a series of different distances (d) from end A.

ANALYSIS AND CALCULATIONS:

Plot a graph of T against d.

From the graph record a series of values of the simple equivalent pendulum (L).

Calculate the value of g from the graph or from the formula:

 $T^2 = 4\pi^2 L/g$



Torsion Pendulum:

1. Introduction

Torsion is a type of stress, which is easier to explain for a uniform wire or a rod when one end of the wire is fixed, and the other end is twisted about the axis of the wire by an external force. The external force causes deformation of the wire and appearance of counterforce in the material. If this end is released, the internal torsion force acts to restore the initial shape and size of the wire. This behavior is similar to the one of the released end of a linear spring with a mass attached.

Attaching a mass to the twisting end of the wire, one can produce a torsion pendulum with circular oscillation of the mass in the plane perpendicular to the axis of the wire.

To derive equations of rotational motion of the torsion pendulum, it would be useful to recall a resemblance of quantities in linear and rotational motion. We know that if initially a mass is motionless, its linear motion is caused by force F; correspondingly, if an extended body does not rotate initially, its rotation is caused by torque τ . The measure of inertia in linear motion is mass, m, while the measure of inertia in rotational motion is the moment of inertia about an axis of rotation, I. For linear and angular displacement in a one-dimensional

problem, we use either x or θ . Thus, the two equations of motion are:

 $F_x = ma_x$ and $\tau = I\alpha$ (1)

where a_x and α are the linear and the angular acceleration.

If the linear motion is caused by elastic, or spring, force, the Hooke's law gives $F_x = -kx$, where k is the spring constant. If the rotation is caused by torsion, the Hooke's law must result in

 $\tau = -\kappa \theta$

where *x* is the torsion constant, or torsional stiffness, that depends on properties of the wire. It is essentially a measure of the amount of torque required to rotate the free end of the wire 1 radian.

Your answer to the Preparatory Question 2 gives the following relationship between the moment of inertial of an oscillating object and the period of oscillation Tas:

This relationship is true for oscillation where damping is negligible and can be ignored. Otherwise the relationship between I and κ is given by

$$I = \frac{\kappa}{\omega_0^2} \tag{3*}$$

(3**)

where ω_0 can be found from $\omega = \sqrt{\omega_0^2 - \left(\frac{c}{2I}\right)^2}$

 $\omega = \frac{2\pi}{T} = 2\pi f; f$ is the frequency of damped oscillation; and c is the *damping coefficient*.

The relationship between the torsion constant κ and the diameter of the wired is given in [3] (check your answer to the Preparatory Question 1) as

$$\kappa = \frac{\pi G d^4}{32l} \tag{4}$$

where l is the length of the wire and G is the shear modulus for the material of the wire.

As any mechanical motion, the torsional oscillation is damped by resistive force originating from excitation of thermal modes of oscillation of atoms inside the crystal lattice of the wire and air resistance to the motion of the oscillating object. We can estimate the torque of the resistive force as being directly proportional to the angular speed of the twisting wire, i.e. the torque $\tau_R = -cd\theta/dt$ (recall the drag force on mass on spring in viscose medium as R = -bv). Combining Eq.(1), (2) and the expression for τ_R , we obtain the equation of motion of a torsional pendulum as follows:

$$I\frac{d^{2}\theta}{dt^{2}} + c\frac{d\theta}{dt}\kappa\theta = 0$$
 (5)

The solution of Eq.(5) is similar to the solution of the equation for damped oscillation of a mass on spring and is given by:

 $\boldsymbol{\theta} = A e^{-\alpha t} \cos(\omega t + \boldsymbol{\phi}) \tag{6}$

where $\alpha = c/2I$

(7)

and $\alpha = \beta^{-1}$ with β being the time constant of the damped oscillation; *c* is the damping coefficient; ω is the angular frequency of torsional oscillation measured in the experiment; and φ can be made zero by releasing the object on the wire at a position of the greatest deviation from equilibrium.

Equation (6) can be used to calculate c (damping coefficient) and β (time constant = amount of time to decaye times) with DataStudio interface and software.

Another important formula is $\alpha = \omega_0/2Q$, where Q is the *quality factor* and $\omega_0^2 = \kappa/I$ (see Eq.3'). The ratio $\zeta = \alpha/\omega_0 = (2Q)^{-1}$ (8)

is called the *damping ratio*.

Free vibration of One Degree of Freedom Systems

Free vibration of a system is vibration due to its own *internal forces* (free of external impressive forces). It is initiated by an initial deviation (an energy input) of the system from its static equilibrium position. Once the initial deviation (a displacement or a velocity or both) is suddenly withdrawn, the strain energy stored in the system forces the system to return to its original, static equilibrium configuration. Due to the inertia of the system, the system will not return to the equilibrium configuration in a straightforward way. Instead it will oscillate about this position — free vibration.

A system experiencing free vibration oscillates at one or more of its natural frequencies, which are properties of its mass and stiffness distribution. If there is no damping (an undamped system), the system vibrates at the (*undamped*) frequency (frequencies) forever. Otherwise, it vibrates at the (*damped*) frequency (frequencies) and dies out gradually. When damping is not large, as in most cases in engineering, undamped and damped frequencies are very close. Therefore usually no distinction is made between the two types of frequencies.

The number of natural frequencies of a system equals to the number of its degrees-of-freedom. Normally, the low frequencies are more important.

Damping always exists in materials. This damping is called material damping, which is always positive (dissipating energy). However, air flow, friction and others may 'present' negative damping. **Undamped Free Vibration**

Equation of motion based on the free-body diagram
$$m\mathbf{x} + k\mathbf{x} = 0$$

 $\mathbb{M} + \omega_{\rm h}^2 x = 0$

 $\omega_{n} = \sqrt{\frac{k}{m}}$

natural frequency

period

$$\tau = 2\pi \sqrt{\frac{m}{k}}$$

 $x(t) = A \sin \psi t + B \cos \psi t$

A and B are determined by the *initial conditions*.

Sin or Cos



Systems involving rotational degrees-of-freedom are always more difficult to deal with, in particular when translational degrees-of-freedom are also present. Gear care is needed to identify both degrees-of-freedom and construct suitable equations of motion.

Damped Free Vibration (first hurdle in studying vibration)





3. critically damped motion () $\zeta = 1$



 $x(t) = (A + Bt) \exp(-\omega t)$

4. negative damping of $\zeta < 0$ as a special case of $\zeta < 1$:



Divergent oscillatory motion (flutter) due to negative damping

Determination of Damping

$$x(t) = X \exp(-t \mu t) \sin(\mu t + \phi)$$



 $2 \exp(-0.05\pi t) \sin(0.9988\pi t + \phi)$

two consecutive peaks:

$$x_{1} = X \exp(-\zeta_{1}t_{1}) \sin(\zeta_{1}t_{1} + \phi)$$

$$x_{2} = X \exp(-\zeta_{1}t_{2}) \sin(\zeta_{1}t_{2} + \phi) = X \exp(-\zeta_{1}t_{2}) \sin(\zeta_{1}t_{1} + \phi)$$

$$\delta = \ln \frac{x_{1}}{2} = \zeta_{0}t_{1} \text{ and } \zeta = \frac{\delta}{2}$$

logarithm decrement \Rightarrow

$$\delta = \ln \frac{x_1}{x_2} = \zeta \omega \tau \quad \text{n d} \quad \zeta = \frac{\delta}{\omega \tau}$$

Example:

The 2nd and 4th peaks of a damped free vibration measured are respectively 0.021 and 0.013. What is damping factor? **Solution**:

$$\frac{x(t_2)}{x(t_4)} = \exp(\psi_n 2_d) \longrightarrow 2\psi_n = \ln\left(\frac{x(t_2)}{x(t_4)}\right)$$

$$2\psi_n = 2\psi_n \frac{2\pi}{\psi_n \sqrt{1-\xi^2}} = \frac{4\pi\xi}{\sqrt{1-\xi^2}} = \ln\left(\frac{x(t_2)}{x(t_2)}\right)$$

$$= \ln\left(\frac{x(t_4)}{x(t_2)}\right)$$

$$= \ln\left(\frac{x(t_4)}{x(t_2)}\right)$$

$$= \ln\left(\frac{x(t_4)}{x(t_2)}\right)$$

$$= \ln\left(\frac{x(t_4)}{x(t_2)}\right)$$

$$= \ln\left(\frac{x(t_4)}{x(t_2)}\right)$$

$$= \ln\left(\frac{x(t_4)}{x(t_2)}\right)$$

If a small damping is assumed, $2\pi = 4\pi = \ln \left| \frac{1}{x(t^{-4})} \right|$. This leads to

$$\zeta = \frac{1}{4\pi} \ln \left(\frac{x(t_2)}{x(t_4)} \right) = 0.0382 = 3.82\%$$

such assumption is made, then an not

If such an assumption is not made, then
$$\frac{\zeta}{\sqrt{1-\zeta^2}} = \frac{1}{4\pi} \ln\left(\frac{x(t_2)}{x(t_4)}\right)$$
$$\frac{\zeta}{1-\zeta^2} = \left| \frac{1}{4\pi} \ln\left(\frac{x(t_2)}{x(t_4)}\right) \right|^2$$
. This leads to
$$\frac{1}{4\pi} \ln\left(\frac{x(t_2)}{x(t_4)}\right) = 0.0381 = 3.81\%$$
$$\sqrt{1 + \left[\frac{1}{4\pi} \ln\left(\frac{x(t_2)}{x(t_4)}\right)\right]^2}$$
. So virtually the same value.

and hence

General differential equations

If

$$a_{n}\frac{d^{n}x}{dt^{n}} + a_{n-1}\frac{d^{n-1}x}{dt^{n-1}} + \dots + a_{1}\frac{dx}{dt^{1}} + a_{0} = 0$$

first solve the characteristic equation

$$a h + a h^{n-1} + \dots + a + a = 0$$

If all roots λ_i are distinct, then the general solution is

$$x(t) = \sum_{j=1}^{n} b_{j} \exp(\lambda_{j} t)$$

where b_j are constants to be determined.

If there are repeated roots, t^m (integer m > 1) appears in a solution. These are not interesting cases for mechanical vibration.

 λ in response to the change of a parameter reveal stability properties